

INVESTIGATION OF THE STRUCTURE OF

INTACT CORE

IN  
FULL CONE SPRAYS

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20 TH DISC MEETING  
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SANDIA LABS

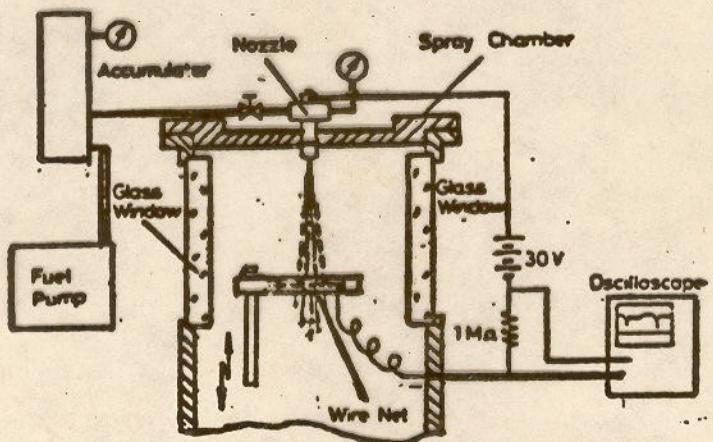


Fig. 8 Experimental apparatus for measuring break-up length

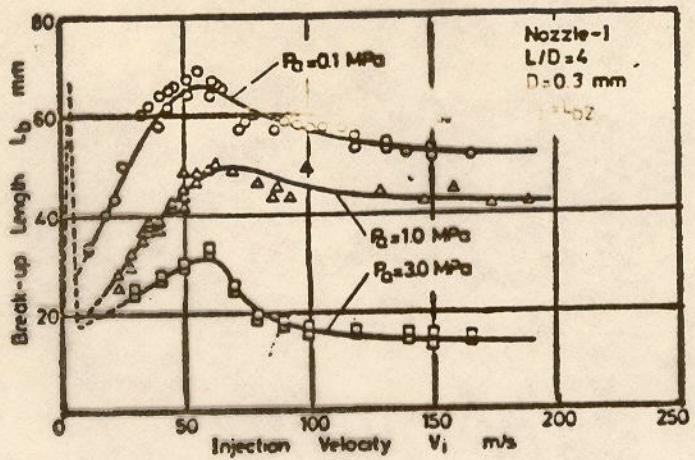
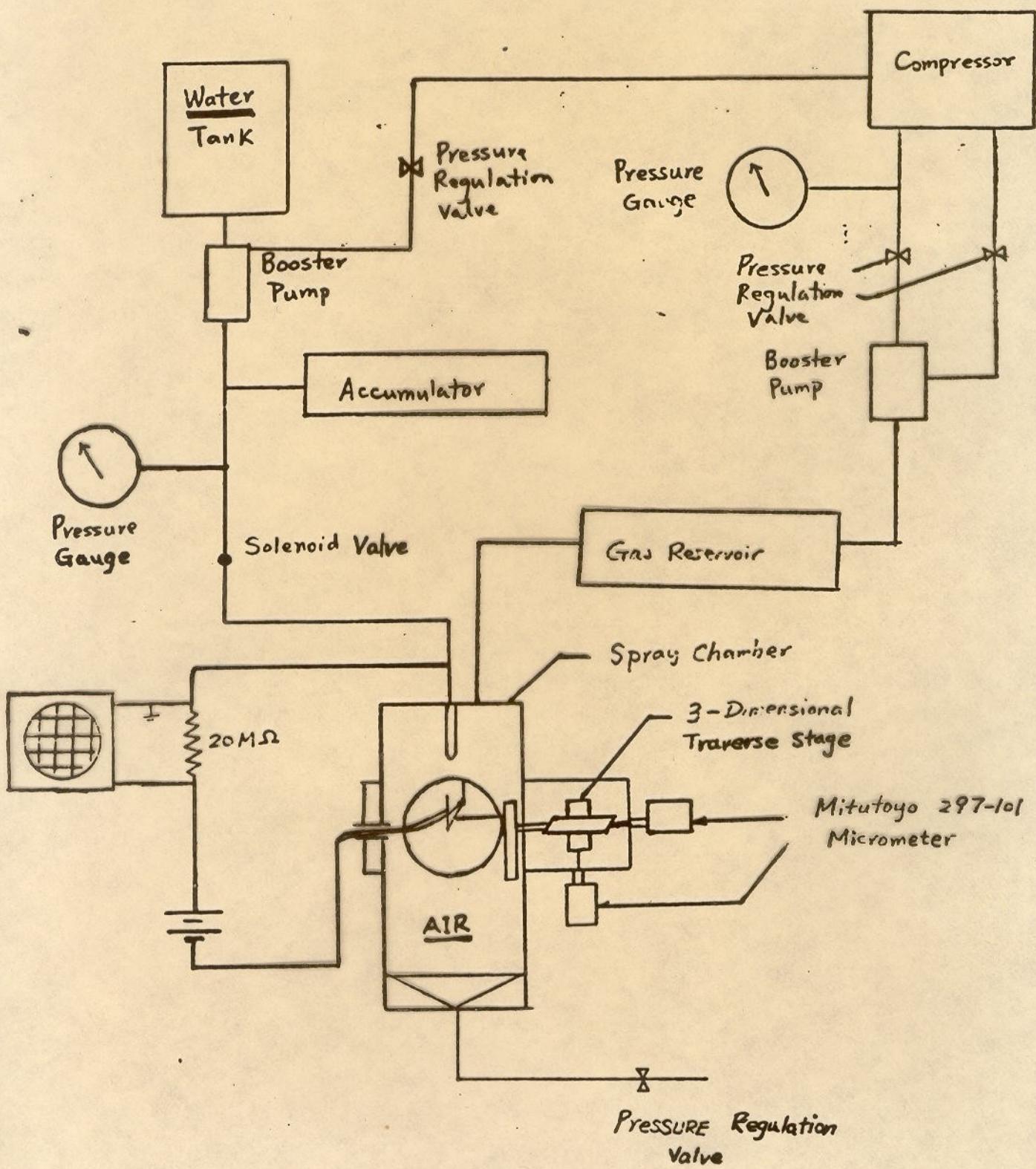


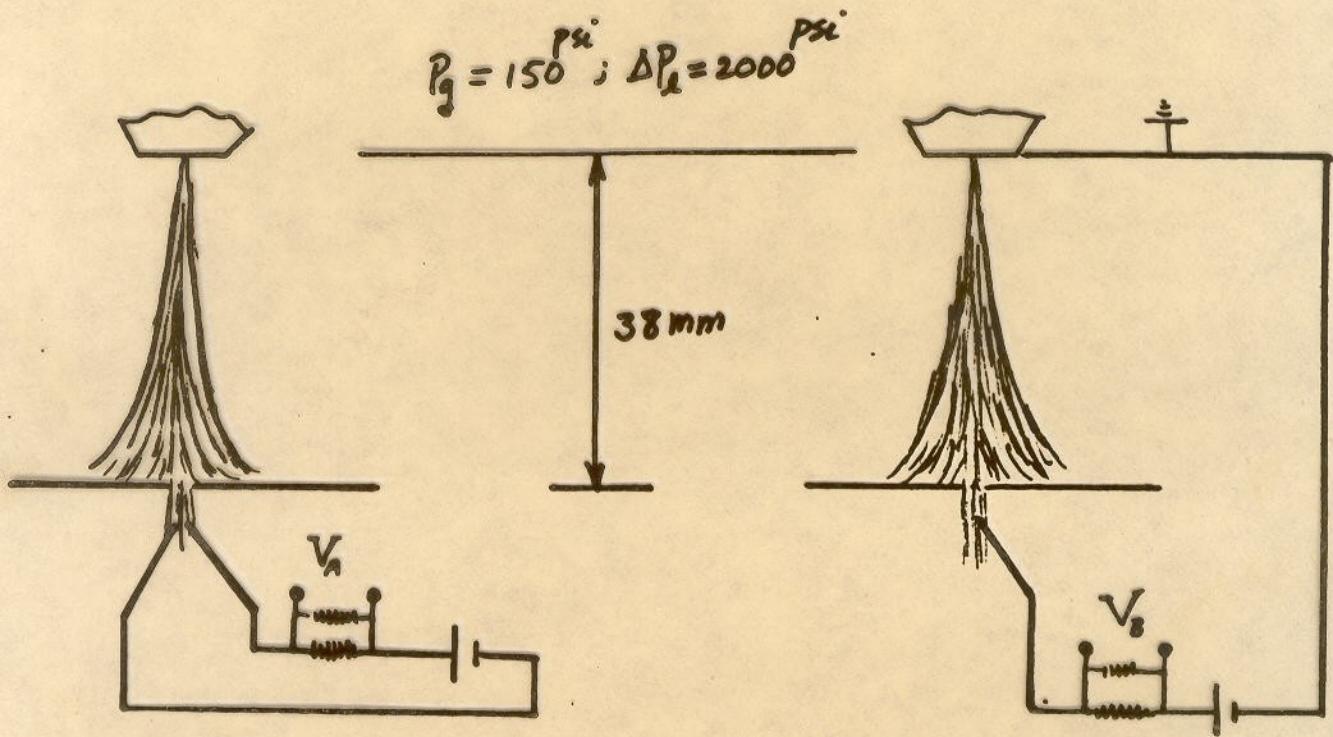
Fig. 11 Break-up length

Arai M., Tabata M., Hiroyasu H., and Shimizu M.  
 "Disintegrating Process and Spray Characterization  
 of Fuel Jet Injected by a Diesel Nozzle"  
 SAE Tech. Paper No. 840275 (1984)

Using water, we decided to check Arai, et al. findings and, most of all, to try and determine the shape of the intact core using a 1 mm diameter bar that can be moved in 3-D instead of the screen that was moved only axially.

Our set up is shown:





$$V_h = 120 - 150 \text{ mV}$$

$$(3331 - 2620 \text{ M}\Omega)$$

$$V_A (\text{when blocked}) = 9-15 \text{ mV}$$

$$\underline{\underline{\text{Gap Size} \approx 250 \mu\text{m}}}$$

$$V_B = 30 - 50 \text{ mV}$$

$$(12,800 - 7,568 \text{ M}\Omega)$$

$$V_B (\text{when blocked}) = 5-18 \text{ mV}$$

Our first important finding is that a current is measured when drops impinge on a bar even if the intact core is not touched.

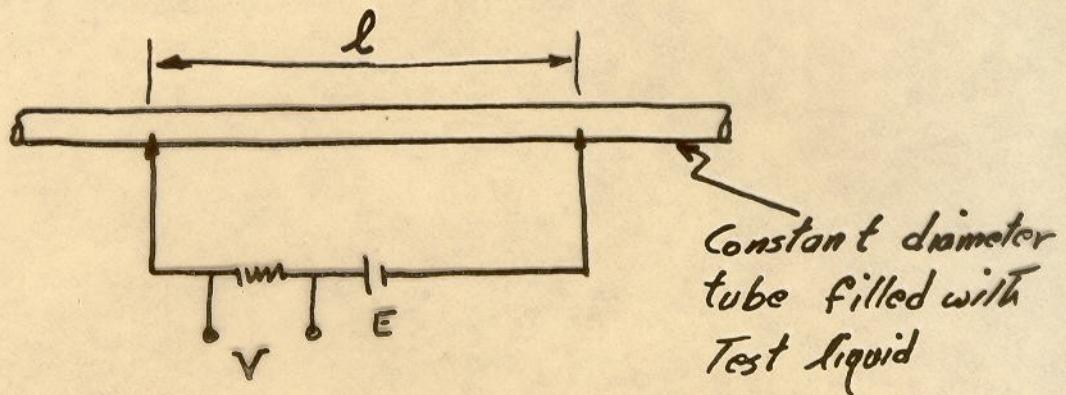
In fact a current is measured also when  
a spray flows between two needles and  
the needles themselves are dry.

Thus the current measurement by Hiroyasu does not imply necessarily the existence of an intact core between the nozzle and the measurement station.

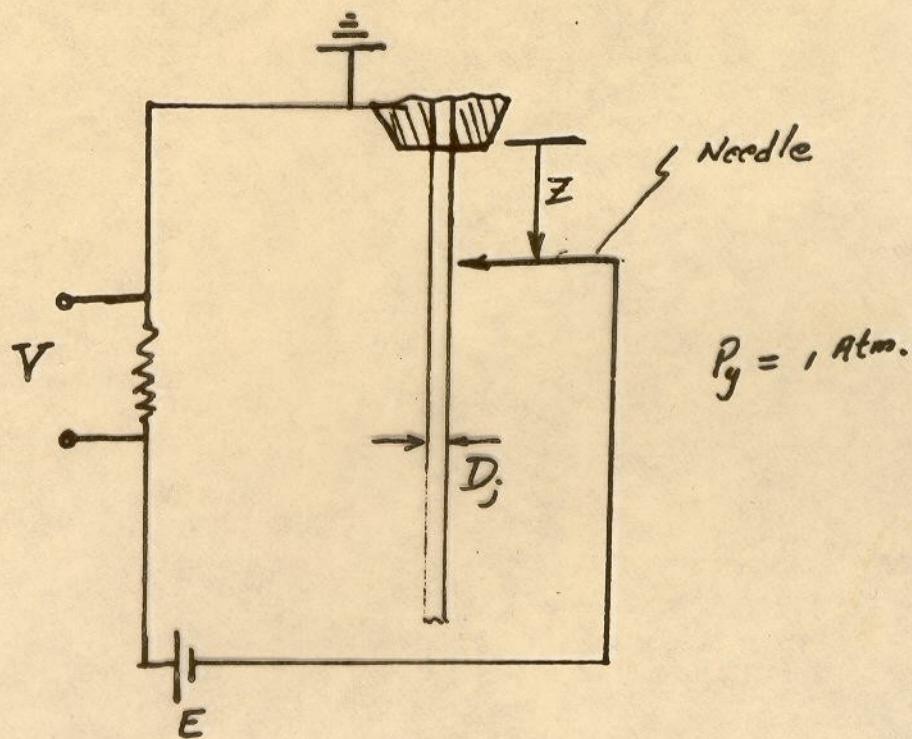
The current carried by the drops, or across the spray, sets a limit also on our own measurements which will be discussed later.

In our efforts to determine the length and the shape of the core, we first measured the resistance of an intact cylindrical column of water to check and calibrate the system.

## Resistivity Measurement



## Low Velocity (Laminar) jet experiment



Diameter of the jet was measured using 26X magnification camera.

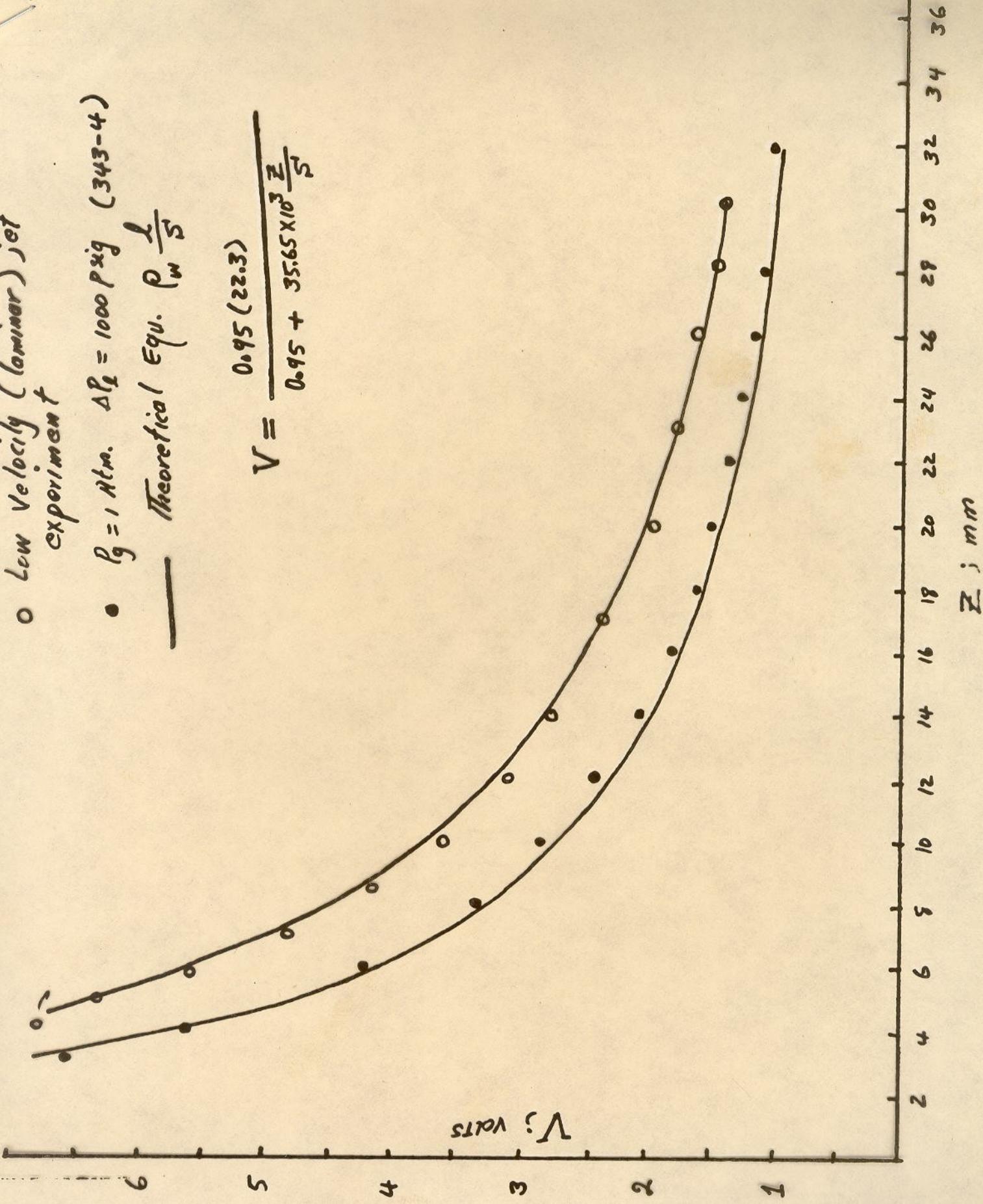
$$D_j \approx 308 \pm 10 \mu\text{m} \quad (D_n = 343 \mu\text{m})$$

○ Low Velocity (laminar) jet  
experiment

●  $\rho_g = 1 \text{ atm}$ .  $\Delta \rho = 1000 \text{ Pug}$  (343-4)  
— Theoretical Equ.  $P_w \frac{L}{S}$

$$V = \frac{0.95(22.3)}{0.95 + 35.65 \times 10^3 \frac{Z}{S}}$$

VOLTS

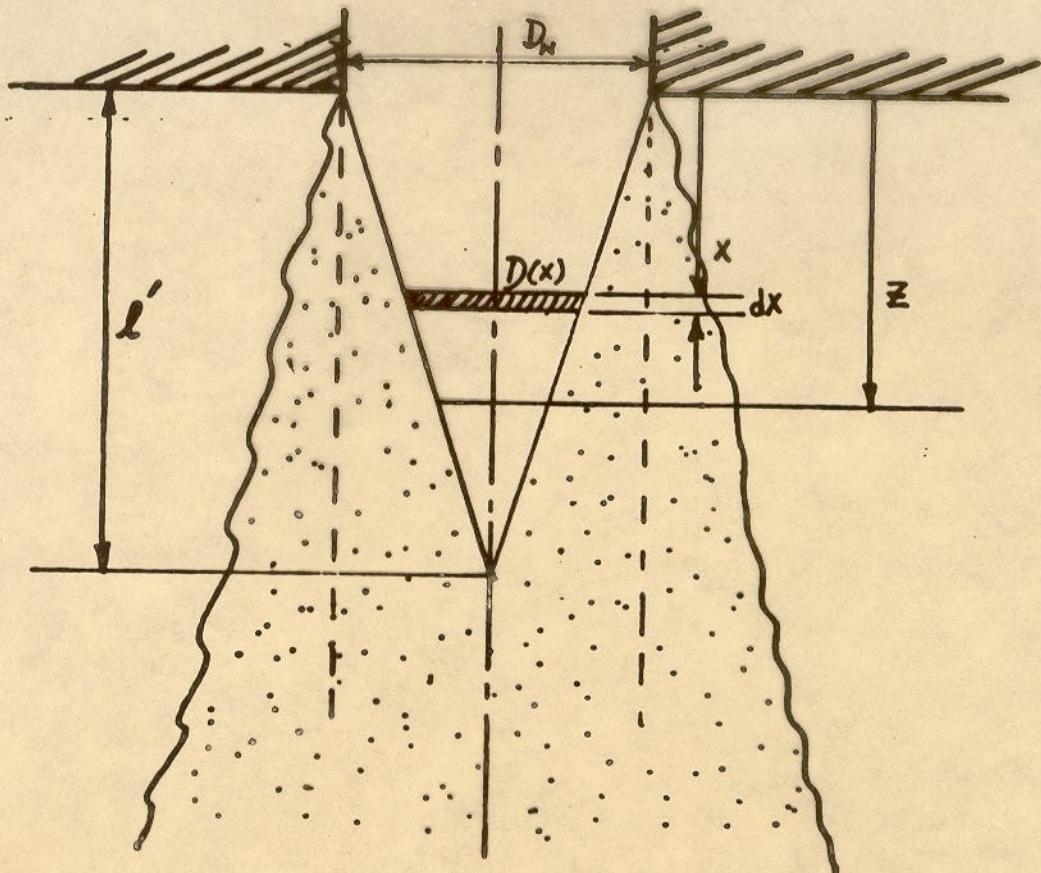


## Conclusion (so far!)

Resistance equ. ( $\rho \frac{l}{s}$ ) for metallic conductors is also applicable for an intact column of water.

## Extension

Since water behaves ... similar to metallic conductor and if there is an intact core of unknown geometry, then a measurement of the resistance of the core at various locations along the jet may lead to its shape.



$$D(x) = D_N f(x)$$

$$dR = \rho \frac{dx}{S} ; \quad S = \frac{\pi D_N^2 f(x)^2}{4}$$

$$dR = \frac{4\rho}{\pi D_N^2} \cdot \frac{dx}{f(x)^2}$$

$$R_{\text{spray}}(z) = \frac{4\rho}{\pi D_N^2} \cdot \int_0^z \frac{dx}{[f(x)]^2}$$

If linear model :  $D(x) = D_N \left(1 - \frac{x}{\ell'}\right)$

$$\alpha \equiv R_{\text{spray}}(z) / (4\rho \ell / \pi D_N^2) \equiv \frac{\ell'}{\ell' - z}$$

Table of Test Conditions

Test case	Nozzle	Dia - 4D (Dia μm)	Liquid	$\Delta P_e$ psig	$P_g$ psig	$\frac{P_e}{P_g} \left( \frac{Re}{We} \right)^2$
R1	X-9	508-4	H <sub>2</sub> O	2000	150	36.6
A2	"	"	"	"	420	14.4
B1	X-8	343-4	"	"	150	36.6
B2	"	"	"	"	420	14.4
C1	X-5	305-4	"	"	150	36.6
C2	"	"	"	"	420	14.4
D1	X-6	127-4	"	"	150	36.6
D2	"	"	"	"	420	14.4

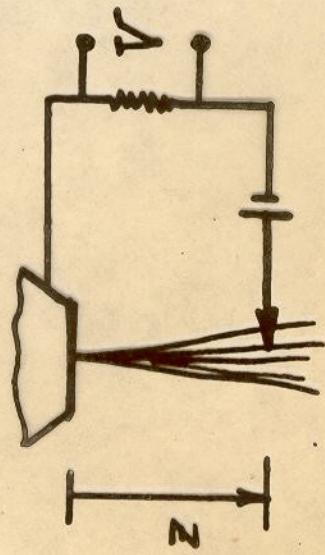
•  $P_{\text{gas}} = 1 \text{ atm}$      $\Delta P_t = 1000 \text{ Pa}$

+

case B1

Δ case B2

$\Delta_f^*$



5

4

3.5

2.5

2

1

-1

0

4

12

16

20

24

28

32

+

+

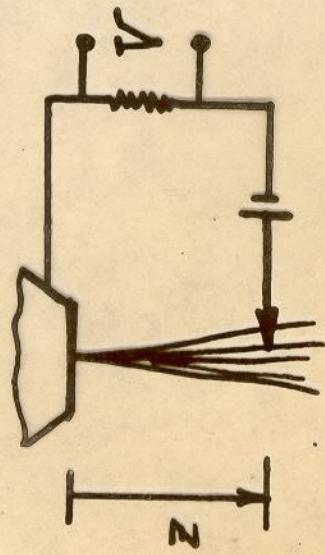
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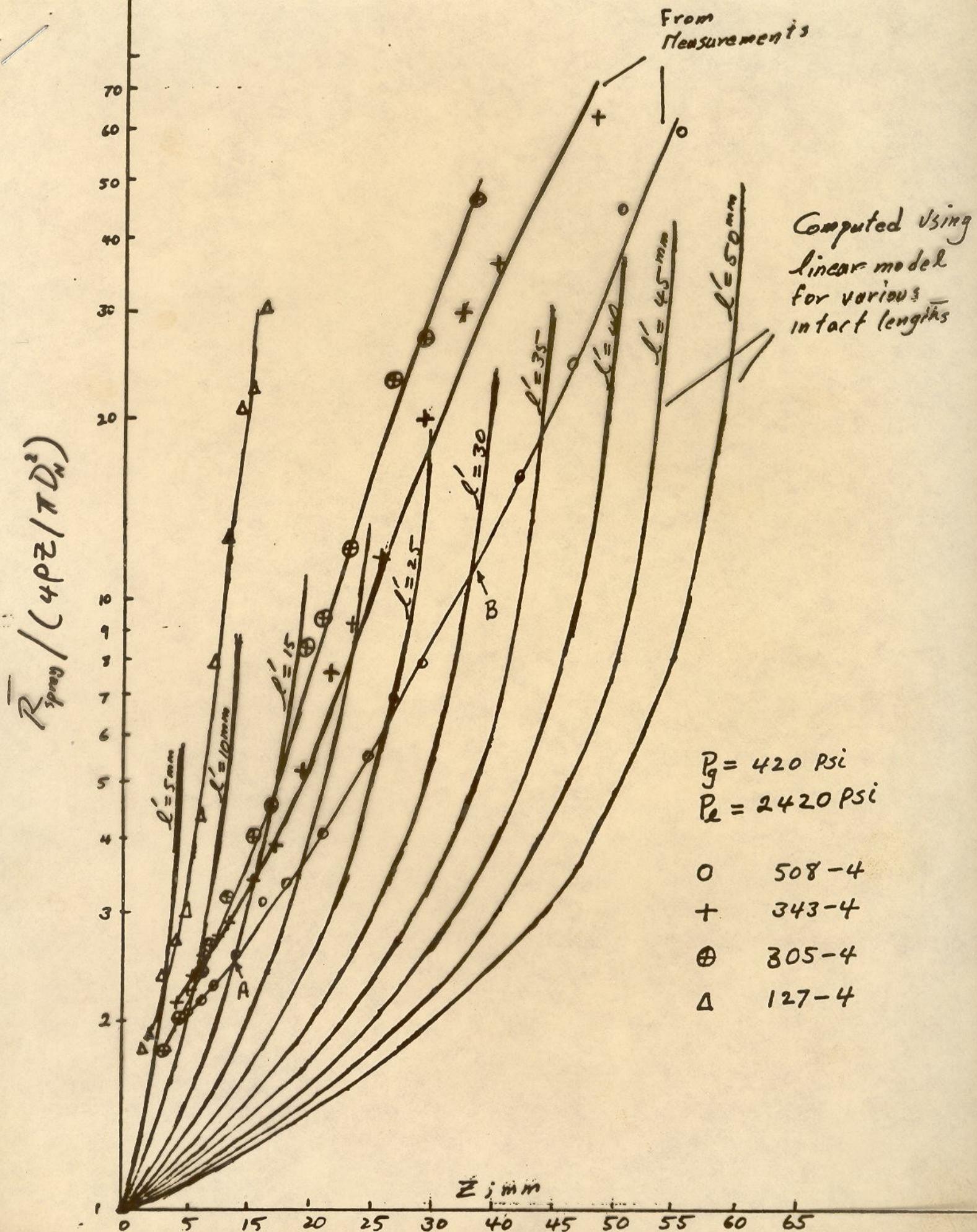
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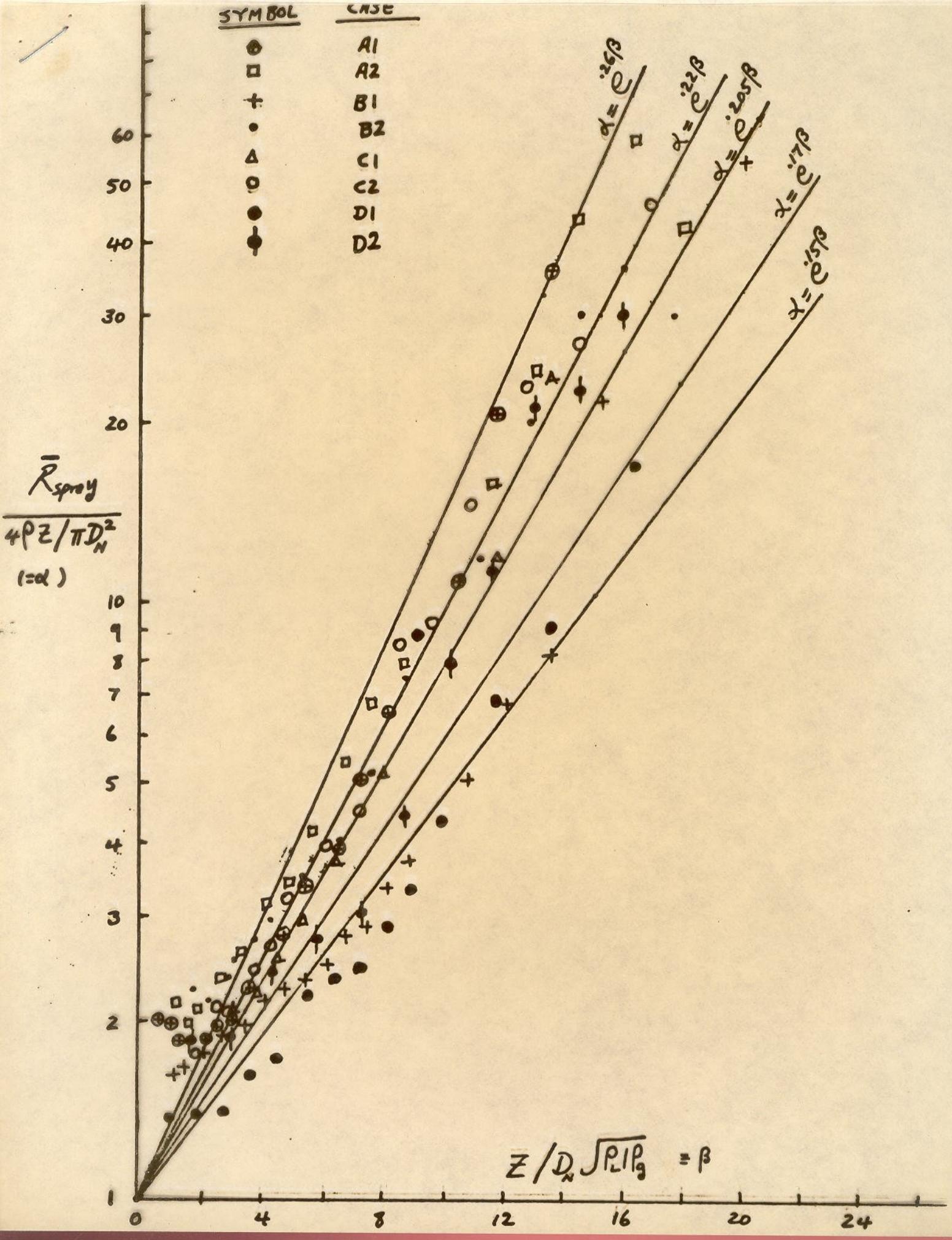
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If it is assumed that the intact core is triangular in shape. Then one can compute its resistance versus axial distance from the nozzle and compare it with the experimental results.

In the next figure, at "A" the resistance of the intact core is equivalent to the resistance of the frustum of a cone with length of  $l' = 15 \text{ mm}$  and similarly at "B" the resistance of the core is equal to the frustum of a cone with  $l' = 30 \text{ mm}$ . This indicates that the shape of the core is convex towards the jet axis.



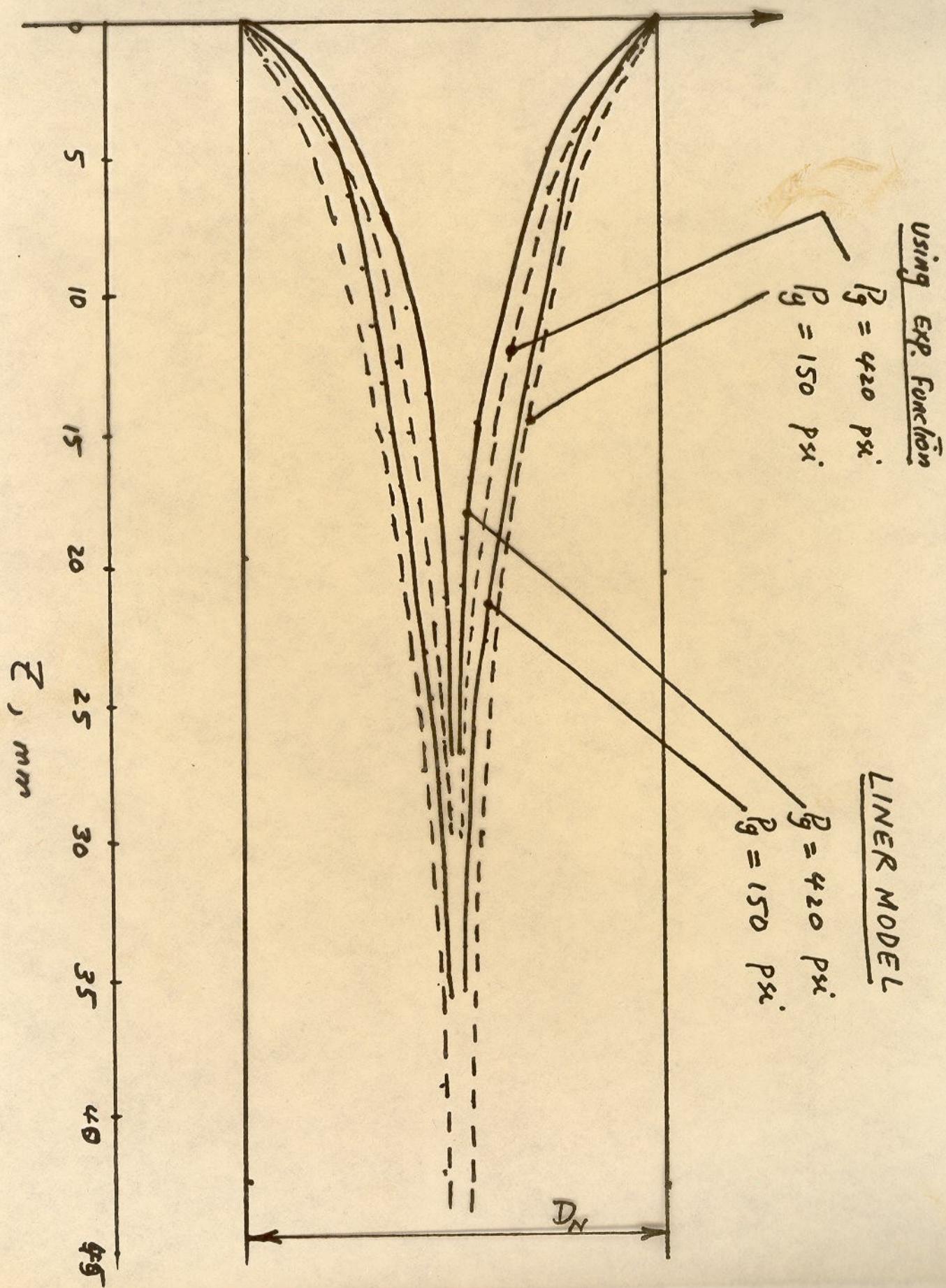


An exponential function has been tried for the experimental data in the form  $\alpha = e^{\alpha \beta}$  from which one can compute the shape of the core. Results are good for higher pressure but not satisfactory at low chamber pressure particularly near the exit of the nozzle.

One reason for selecting an exponential function of this form is that in the limit as  $\beta \rightarrow 0$  (close to the exit of the nozzle)  $D(x)/D_N \rightarrow 1 - z\alpha\sqrt{P_g/P_L}$ . Such a behavior is predicted by the supplemented - aerodynamic theory of the jet breakup in the vicinity of the nozzle exit.

Using  $\alpha = e^{\alpha \beta}$  one could derive an equation for the shape of the core as:

$$\frac{D(x)}{D_N} = f(x) = \frac{\exp\left[-\frac{\alpha z}{2D_N\sqrt{P_L/P_g}}\right]}{\left[1 + \frac{\alpha z}{D_N\sqrt{P_L/P_g}}\right]^{1/2}}$$



## CONCLUSIONS

1. Measurements were made of the shape of the intact core in atomizing sprays using a 1mm diameter bar. They indicate that the mean core boundary is convex towards the jet axis. The length of the core increases with nozzle diameter and liquid-to-gas density ratio.
2. The technique used tends to overestimate the length of the core due to the current carried by the charged drops and by the gas-drop mixture.
3. Measurement of the intact length geometry using a sharp needle to locate the core boundary is feasible but only up to 10 to 15 nozzle diameter from the nozzle where the spray interference is negligible.